

THREE-DIMENSIONAL C_{12} -MANIFOLDS

GHERICI BELDJILALI

ABSTRACT. The present paper is devoted to three-dimensional C_{12} -manifolds (defined by D. Chinea and C. Gonzalez), which are never normal. We study their fundamental properties and give concrete examples. As an application, we study such structures on three-dimensional Lie groups.

1. INTRODUCTION

In [6], D. Chinea and C. Gonzalez obtained a classification of the almost contact metric manifolds, studying the space that possess the same symmetries as the covariant derivative of the fundamental 2-form. This space is decomposed into twelve irreducible components C_1, \dots, C_{12} . In dimension 3, the classes C_i reduce to the following classes: $|C|$ class of cosymplectic manifolds, C_5 class of β -Kenmotsu manifolds, C_6 class of α -Sasakian manifolds, C_9 -manifolds and C_{12} -manifolds.

Most of the research related to almost contact metric structures is concerned with the normal structures which contain the first three classes. Regarding the C_{12} class which is not normal, only two papers address this subject. In the first one [5], the authors developed a systematic study of the curvature of the Chinea–Gonzalez class $C_5 \oplus C_{12}$ and obtain some classification theorems for those manifolds that satisfy suitable curvature conditions. This class is defined by using a certain function α and when this function vanishes the class $C_5 \oplus C_{12}$ reduces to class C_{12} . The second paper [3] contains new results on a particular three-dimensional C_{12} -manifolds with a class of concrete illustrative examples.

The present paper is devoted to three-dimensional C_{12} -manifolds. We present a detailed study of such class in dimension three and we construct a class of examples. As an application, we give all C_{12} -structures on Lie algebras of dimension 3.

First of all, we will start by introducing the basic concepts that we need in this research.

2. ALMOST CONTACT MANIFOLDS

An odd-dimensional Riemannian manifold (M^{2n+1}, g) is said to be an *almost contact metric manifold* if there exist on M a $(1, 1)$ -tensor field φ , a vector field ξ

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(called the structure vector field) and a 1-form η such that

$$\begin{cases} \eta(\xi) = 1, \\ \varphi^2(X) = -X + \eta(X)\xi, \\ g(\varphi X, \varphi Y) = g(X, Y) - \eta(X)\eta(Y) \end{cases}$$

for any vector fields X, Y on M .

In particular, in an almost contact metric manifold we also have

$$\varphi\xi = 0 \quad \text{and} \quad \eta \circ \varphi = 0.$$

The fundamental 2-form ϕ is defined by

$$\phi(X, Y) = g(X, \varphi Y).$$

It is known that the almost contact structure (φ, ξ, η) is said to be normal if and only if

$$N^{(1)}(X, Y) = N_\varphi(X, Y) + 2d\eta(X, Y)\xi = 0$$

for any X, Y on M , where N_φ denotes the Nijenhuis torsion of φ , given by

$$N_\varphi(X, Y) = \varphi^2[X, Y] + [\varphi X, \varphi Y] - \varphi[\varphi X, Y] - \varphi[X, \varphi Y].$$

Given an almost contact structure, one can associate in a natural manner an almost CR-structure $(\mathcal{D}, \varphi|_{\mathcal{D}})$, where $\mathcal{D} := \text{Ker}(\eta) = \text{Im}(\varphi)$ is the distribution of rank $2n$ transversal to the characteristic vector field ξ . If this almost CR-structure is integrable (i.e., $N_\varphi = 0$) the manifold M^{2n+1} is said to be CR-integrable. It is known that normal almost contact manifolds are CR-manifolds.

For more background on almost contact metric manifolds, we recommend the references [1, 4, 9].

3. THREE-DIMENSIONAL C_{12} -MANIFOLDS

In the classification of D. Chinea and C. Gonzalez [6] of almost contact metric manifolds there is a class called C_{12} -manifolds which can be integrable but never normal. In this classification, C_{12} -manifolds are defined by

$$(\nabla_X \phi)(Y, Z) = \eta(X)\eta(Z)(\nabla_\xi \eta)\varphi Y - \eta(X)\eta(Y)(\nabla_\xi \eta)\varphi Z.$$

In [3] and [5], the $(2n + 1)$ -dimensional C_{12} -manifolds are characterized by

$$(\nabla_X \varphi)Y = \eta(X)(\omega(\varphi Y)\xi + \eta(Y)\varphi\psi) \quad (3.1)$$

for any X and Y vector fields on M , where $\omega = -(\nabla_\xi \xi)^b = -\nabla_\xi \eta$ and ψ is the vector field given by

$$\omega(X) = g(X, \psi) = -g(X, \nabla_\xi \xi)$$

for all X vector field on M .

Moreover, in [3] the $(2n + 1)$ -dimensional C_{12} -manifolds are also characterized by

$$d\eta = \omega \wedge \eta, \quad d\phi = 0 \quad \text{and} \quad N_\varphi = 0.$$

Here, we emphasize that the almost C_{12} -manifolds are defined as follows.

Definition 3.1. Let $(M^{2n+1}, \varphi, \xi, \eta, g)$ be an almost contact metric manifold. M is called *almost C_{12} -manifold* if there exists a closed one-form ω which satisfies

$$d\eta = \omega \wedge \eta \quad \text{and} \quad d\phi = 0.$$

In addition, if $N_\varphi = 0$ we say that M is a C_{12} -manifold.

On the other hand, in [7] the author proved that, for an arbitrary 3-dimensional almost contact metric manifold $(M^3, \varphi, \xi, \eta, g)$, we have

$$\begin{cases} (1) & (\nabla_X \varphi)Y = g(\varphi \nabla_X \xi, Y)\xi - \eta(Y)\varphi \nabla_X \xi, \\ (2) & d\phi = (\operatorname{div} \xi)\eta \wedge \phi, \\ (3) & d\eta = \eta \wedge (\nabla \xi \eta) + \frac{1}{2}(\operatorname{tr}_g(\varphi \nabla \xi))\phi. \end{cases}$$

Then, for any 3-dimensional almost C_{12} -manifold $(M^3, \varphi, \xi, \eta, g)$ we get

$$(\nabla_X \varphi)Y = g(\varphi \nabla_X \xi, Y)\xi - \eta(Y)\varphi \nabla_X \xi \tag{3.2}$$

and

$$\operatorname{div} \xi = \operatorname{tr}_g(\varphi \nabla \xi) = 0.$$

Now we shall introduce another possible sufficient and necessary condition of the integrability of almost C_{12} -manifolds.

Proposition 3.2. *The almost C_{12} -structure (φ, ξ, η, g) is integrable if and only if, for all X and Y vector fields on M , we have*

$$(\nabla_{\varphi X} \varphi)Y - \varphi(\nabla_X \varphi)Y = -g(\nabla_X \xi, Y)\xi - \eta(X)(\omega(Y)\xi - \eta(Y)\psi). \tag{3.3}$$

Proof. We know that

$$N_\varphi(X, Y) = (\varphi \nabla_Y \varphi - \nabla_{\varphi Y} \varphi)X - (\varphi \nabla_X \varphi - \nabla_{\varphi X} \varphi)Y.$$

Suppose that $N_\varphi = 0$ and put

$$\begin{aligned} T(X, Y, Z) &= g(\varphi(\nabla_X \varphi)Y - (\nabla_{\varphi X} \varphi)Y, Z) \\ &= -g((\nabla_X \varphi)Y, \varphi Z) - g((\nabla_{\varphi X} \varphi)Y, Z). \end{aligned}$$

One can easily get

$$T(X, Y, Z) = T(Y, X, Z). \tag{3.4}$$

On the other hand, using formulas

$$\nabla_X(\varphi Y) = (\nabla_X \varphi)Y + \varphi \nabla_X Y \quad \text{and} \quad g(\varphi X, Y) = -g(X, \varphi Y),$$

we can get

$$g((\nabla_X \varphi)Y, Z) = -g(Y, (\nabla_X \varphi)Z),$$

and by straightforward computation we have

$$T(X, Y, Z) = -T(X, Z, Y) + g(\nabla_X \xi, Y)\eta(Z) + g(\nabla_X \xi, Z)\eta(Y). \tag{3.5}$$

Now, using formulas (3.4) and (3.5) we obtain

$$\begin{aligned}
 T(X, Y, Z) &= T(Y, X, Z) \\
 &= -T(Y, Z, X) + g(\nabla_Y \xi, X)\eta(Z) + g(\nabla_Y \xi, Z)\eta(X) \\
 &= T(Z, X, Y) - g(\nabla_Z \xi, X)\eta(Y) - g(\nabla_Z \xi, Y)\eta(X) \\
 &\quad + g(\nabla_Y \xi, X)\eta(Z) + g(\nabla_Y \xi, Z)\eta(X) \\
 &= -T(X, Y, Z) + g(\nabla_X \xi, Y)\eta(Z) + g(\nabla_X \xi, Z)\eta(Y) \\
 &\quad - g(\nabla_Z \xi, X)\eta(Y) - g(\nabla_Z \xi, Y)\eta(X) \\
 &\quad + g(\nabla_Y \xi, X)\eta(Z) + g(\nabla_Y \xi, Z)\eta(X),
 \end{aligned}$$

which implies

$$2T(X, Y, Z) = (g(\nabla_X \xi, Y) + g(\nabla_Y \xi, X))\eta(Z) + 2d\eta(X, Z)\eta(Y) + 2d\eta(Y, Z)\eta(X).$$

Since the structure is almost C_{12} -structure, we have

$$\begin{aligned}
 2d\eta(X, Y) &= g(\nabla_X \xi, Y) - g(\nabla_Y \xi, X) \\
 &= \omega(X)\eta(Y) - \eta(X)\omega(Y),
 \end{aligned}$$

therefore

$$T(X, Y, Z) = g(\nabla_X \xi, Y)\eta(Z) + \eta(X)(\omega(Y)\eta(Z) - \eta(Y)\omega(Z)),$$

which gives our formula (3.3). The proof of the converse is direct. \square

We summarize all the above in the following main theorem.

Theorem 3.3. *Let $(M^3, \varphi, \xi, \eta, g)$ be a 3-dimensional almost contact metric manifold. M is a C_{12} -manifold if and only if*

$$\nabla_X \xi = -\eta(X)\psi,$$

where $\psi = -\nabla_\xi \xi$.

Proof. Suppose that $\nabla_X \xi = -\eta(X)\psi$ for all X vector field on M . From (3.2), we get

$$(\nabla_X \varphi)Y = \eta(X)(\omega(\varphi Y)\xi + \eta(Y)\varphi\psi),$$

with $\omega(X) = g(\psi, X)$.

Conversely, assuming that $(M^3, \varphi, \xi, \eta, g)$ is a C_{12} -manifold, this is equivalent to

$$(\nabla_X \varphi)Y = \eta(X)(\omega(\varphi Y)\xi + \eta(Y)\varphi\psi).$$

Setting $Y = \xi$ gives

$$-\varphi\nabla_X \xi = \eta(X)\varphi\psi,$$

and hence

$$\nabla_X \xi = \eta(X)\varphi^2\psi = -\eta(X)\psi. \quad \square$$

The following proposition provides another characterization of 3-dimensional C_{12} -manifolds.

Proposition 3.4. *Let $(M^3, \varphi, \xi, \eta, g)$ be a 3-dimensional almost contact metric manifold. M is a C_{12} -manifold if and only if*

$$\nabla_{\varphi X} \xi = 0.$$

Proof. It is sufficient to prove that $\nabla_{\varphi X} \xi = 0$ and $\nabla_X \xi = -\eta(X)\psi$ are equivalent with $\psi = -\nabla_\xi \xi$. Suppose that $\nabla_X \xi = -\eta(X)\psi$, so it is easy to see that $\nabla_{\varphi X} \xi = 0$.

Conversely, suppose that $\nabla_{\varphi X} \xi = 0$ and replacing X by φX using the formula $\varphi^2 X = -X + \eta(X)\xi$, we obtain $\nabla_X \xi = \eta(X)\nabla_\xi \xi$. This completes the proof. \square

In [3], the authors studied the 3-dimensional unit C_{12} -manifold, i.e. the case where ψ is a unit vector field. We will deal here with the general case, i.e. ψ is not necessarily unitary. For that, taking $V = e^{-\rho}\psi$ where $e^\rho = |\psi|$, we get immediately that $\{\xi, V, \varphi V\}$ is an orthonormal frame. We refer to this basis as *fundamental basis*.

Using this frame, one can get the following:

Proposition 3.5. *For any C_{12} -manifold, for all vector fields X on M we have*

- (1) $\nabla_X \xi = -e^\rho \eta(X)V,$
- (2) $\nabla_\xi V = e^\rho \xi,$
- (3) $\nabla_V V = \varphi V(\rho)\varphi V,$
- (4) $\nabla_\xi \varphi V = 0,$
- (5) $\nabla_V \varphi V = -\varphi V(\rho)V.$

Proof. For the first, using (3.1) for $Y = \xi$ we get

$$(\nabla_X \varphi)\xi = \eta(X)\varphi\psi = e^\rho \eta(X)\varphi V;$$

knowing that $(\nabla_X \varphi)Y = \nabla_X \varphi Y - \varphi \nabla_X Y$ and applying φ we obtain

$$\nabla_X \xi = e^\rho \eta(X)\varphi^2 V = -e^\rho \eta(X)V.$$

For the second, we have

$$2d\omega(\xi, X) = 0 \Leftrightarrow g(\nabla_\xi \psi, X) = g(\nabla_X \psi, \xi) = -g(\psi, \nabla_X \xi) = e^{2\rho} \eta(X),$$

which gives

$$\nabla_\xi \psi = e^{2\rho} \xi \tag{3.6}$$

and then

$$\nabla_\xi V = \nabla_\xi(e^{-\rho}\psi) = -\xi(\rho)V + e^\rho \xi.$$

On the other hand, we have

$$\xi(\rho) = \frac{1}{2}e^{-2\rho}\xi(e^{2\rho}) = \frac{1}{2}e^{-2\rho}\xi(g(\psi, \psi)) = e^{-2\rho}g(\nabla_\xi \psi, \psi) = 0,$$

because of (3.6). Then,

$$\nabla_\xi V = e^\rho \xi.$$

For $\nabla_V V$, we have

$$2d\omega(\psi, X) = 0 \Leftrightarrow g(\nabla_\psi \psi, X) = g(\nabla_X \psi, \psi) = \frac{1}{2}Xg(\psi, \psi) = e^{2\rho}g(\text{grad } \rho, X),$$

i.e. $\nabla_\psi \psi = e^{2\rho} \text{grad } \rho$, which gives $\nabla_V V = \text{grad } \rho - V(\rho)V$.

Also, we have

$$\operatorname{grad} \rho = \xi(\rho)\xi + V(\rho)V + \varphi V(\rho)\varphi V = V(\rho)V + \varphi V(\rho)\varphi V;$$

then,

$$\nabla_V V = \varphi V(\rho)\varphi V.$$

For the rest, just use the formula $\nabla_X \varphi Y = (\nabla_X \varphi)Y + \varphi \nabla_X Y$ noting that

$$(\nabla_V \varphi)X = (\nabla_{\varphi V} \varphi)X = 0. \quad \square$$

It remains to calculate $\nabla_{\varphi V} V$ and $\nabla_{\varphi V} \varphi V$. For that, we have the following lemma.

Lemma 3.6. *For any 3-dimensional C_{12} -manifold, we have*

- (1) $\nabla_{\varphi V} V = (-e^\rho + \operatorname{div} V)\varphi V,$
- (2) $\nabla_{\varphi V} \varphi V = (e^\rho - \operatorname{div} V)V.$

Proof. Since $\{\xi, V, \varphi V\}$ is an orthonormal frame,

$$\nabla_{\varphi V} V = a\xi + bV + c\varphi V.$$

Using Proposition 3.5, we have

$$a = g(\nabla_{\varphi V} V, \xi) = -g(V, \nabla_{\varphi V} \xi) = 0$$

and $b = g(\nabla_{\varphi V} V, V) = 0$. To get the component c , we have

$$\begin{aligned} \operatorname{div} V &= g(\nabla_\xi V, \xi) + g(\nabla_{\varphi V} \psi, \varphi V) \\ &= e^\rho + g(\nabla_{\varphi \psi} \psi, \varphi \psi) \Leftrightarrow g(\nabla_{\varphi V} V, \varphi V) = -e^\rho + \operatorname{div} V; \end{aligned}$$

then,

$$\nabla_{\varphi V} V = (-e^\rho + \operatorname{div} V)\varphi V.$$

Applying φ with (3.1), we obtain

$$\nabla_{\varphi V} \varphi V = (e^\rho - \operatorname{div} V)V. \quad \square$$

According to Proposition 3.5 and Lemma 3.6, the 3-dimensional C_{12} -manifold is completely controllable. That is:

Corollary 3.7. *For any C_{12} -manifold, we have*

$$\begin{aligned} \nabla_\xi \xi &= -e^\rho V, & \nabla_\xi V &= e^\rho \xi, & \nabla_\xi \varphi V &= 0, \\ \nabla_V \xi &= 0, & \nabla_V V &= \varphi V(\rho)\varphi V, & \nabla_V \varphi V &= -\varphi V(\rho)V, \\ \nabla_{\varphi V} \xi &= 0, & \nabla_{\varphi V} V &= (-e^\rho + \operatorname{div} V)\varphi V, & \nabla_{\varphi V} \varphi V &= (e^\rho - \operatorname{div} V)V. \end{aligned}$$

To clarify these notions, we give the following class of examples.

Example 3.8. We denote the Cartesian coordinates in a 3-dimensional Euclidean space $M = \mathbb{R}^3$ by (x, y, z) and define a symmetric tensor field g by

$$g = e^{2f} \begin{pmatrix} \alpha^2 + \beta^2 & 0 & -\beta \\ 0 & \alpha^2 & 0 \\ -\beta & 0 & 1 \end{pmatrix},$$

where $f = f(y) \neq \text{const}$, $\beta = \beta(x)$ and $\alpha = \alpha(x, y) \neq 0$ everywhere are functions on \mathbb{R}^3 with $f' = \frac{\partial f}{\partial y}$. Further, we define an almost contact metric (φ, ξ, η) on \mathbb{R}^3 by

$$\varphi = \begin{pmatrix} 0 & -1 & 0 \\ 1 & 0 & 0 \\ 0 & -\beta & 0 \end{pmatrix}, \quad \xi = e^{-f} \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix}, \quad \eta = e^f(-\beta, 0, 1).$$

The fundamental 1-form η and the 2-form ϕ have the forms

$$\eta = e^f(dx - \beta dy) \quad \text{and} \quad \phi = -2\alpha^2 e^{2f} dx \wedge dy,$$

and hence

$$d\eta = f' e^f (\beta dx \wedge dy + dy \wedge dz) = f' dy \wedge \eta, \\ d\phi = 0.$$

By a direct computation the nontrivial components of $N_{kj}^{(1) i}$ are given by

$$N_{12}^{(1) 3} = \beta f', \quad N_{23}^{(1) 3} = f' \neq 0.$$

But, for all $i, j, k \in \{1, 2, 3\}$,

$$(N_\varphi)_{kj}^i = 0,$$

implying that (φ, ξ, η) becomes integrable non-normal. We have $\omega = f' dy$, i.e. $d\omega = 0$ and knowing that ω is the g -dual of ψ , i.e. $\omega(X) = g(X, \psi)$, we have immediately that

$$\psi = \frac{f'}{\alpha^2} e^{-2f} \frac{\partial}{\partial y}.$$

Thus, (φ, ξ, η, g) is a 3-parameter family of C_{12} structure on \mathbb{R}^3 .

Notice that

$$|\psi|^2 = \omega(\psi) = g(\psi, \psi) = \frac{f'^2}{\alpha^2} e^{-2f}$$

implies that $V = \frac{e^{-f}}{\alpha} \frac{\partial}{\partial y}$ is a unit vector field; then

$$\left\{ \xi = e^{-f} \frac{\partial}{\partial z}, V = \frac{e^{-f}}{\alpha} \frac{\partial}{\partial y}, \varphi V = \frac{e^{-f}}{\alpha} \left(\frac{\partial}{\partial x} + \beta \frac{\partial}{\partial z} \right) \right\}$$

form an orthonormal basis. To verify the result in formula (3.1), the components of the Levi-Civita connection corresponding to g are given by

$$\begin{aligned} \nabla_\xi \xi &= -\frac{f' e^{-f}}{\alpha} V, & \nabla_\xi V &= \frac{f' e^{-f}}{\alpha} \xi, & \nabla_\xi \varphi V &= 0, \\ \nabla_V \xi &= 0, & \nabla_V V &= -\frac{e^{-f}}{\alpha^2} \alpha_1 \varphi V, & \nabla_V \varphi V &= -\varphi \nabla_V V, \\ \nabla_{\varphi V} \xi &= 0, & \nabla_{\varphi V} V &= \frac{e^{-f}}{\alpha^2} (f' \alpha + \alpha_2) \varphi V, & \nabla_{\varphi V} \varphi V &= \varphi \nabla_{\varphi V} V, \end{aligned}$$

where $\alpha_i = \frac{\partial \alpha}{\partial x_i}$. Then, one can easily check that, for all $i, j \in \{1, 2, 3\}$,

$$(\nabla_{e_i} \varphi) e_j = \nabla_{e_i} \varphi e_j - \varphi \nabla_{e_i} e_j = \eta(e_i) (\omega(\varphi e_j) \xi + \eta(e_j) \varphi \psi).$$

Now, we denote by R the curvature tensor and by S the Ricci curvature. From [5, Corollary 3.1], one can get the following:

Corollary 3.9. *For any 3-dimensional C_{12} -manifold, we have*

$$\begin{aligned} R(X, Y)\xi &= -2d\eta(X, Y)\psi - \eta(Y)\nabla_X\psi + \eta(X)\nabla_Y\psi, \\ R(X, \xi)Y &= \omega(X)(\omega(Y)\xi - \eta(Y)\psi) + g(\nabla_X\psi, Y)\xi - \eta(Y)\nabla_X\psi, \\ S(X, \xi) &= -\eta(X)\operatorname{div}\psi. \end{aligned} \tag{3.7}$$

By use of (3.7), we have

$$R(\xi, \psi)\xi = -\omega(\psi)\psi - \nabla_\psi\psi.$$

Therefore

$$g(R(\xi, \psi)\psi, \xi) = -\omega(\psi)^2 - g(\nabla_\psi\psi, \psi).$$

Thus we have

Proposition 3.10. *On 3-dimensional C_{12} -manifolds, the sectional curvature of the plane section spanned by $\{\xi, \psi\}$ is $-\omega(\psi)^2 - g(\nabla_\psi\psi, \psi)$ and if ψ is unitary the sectional curvature is -1 .*

Recall that the conformal curvature tensor vanishes in a 3-dimensional Riemannian manifold, therefore we get (see [2])

$$\begin{aligned} R(X, Y)Z &= g(Y, Z)QX - g(X, Z)QY + S(Y, Z)X - S(X, Z)Y \\ &\quad - \frac{r}{2}(g(Y, Z)X - g(X, Z)Y), \end{aligned} \tag{3.8}$$

where r is the scalar curvature. In the following theorem, we obtain an expression for the Ricci operator in a 3-dimensional C_{12} -manifold.

Theorem 3.11. *In a 3-dimensional C_{12} -manifold, the Ricci operator is given by*

$$QX = (\operatorname{div}\psi)X + (e^\rho - 2\operatorname{div}\psi)\eta(X)\xi - \omega(X)\psi - \nabla_X\psi - \frac{r}{2}\varphi^2X, \tag{3.9}$$

where Q is the Ricci operator defined by

$$S(X, Y) = g(QX, Y). \tag{3.10}$$

Proof. For a 3-dimensional C_{12} -manifold, from (3.7) and (3.8) we have

$$R(X, \xi)\xi = QX + (\operatorname{div}\psi)X - 2(\operatorname{div}\psi)\eta(X)\xi + \frac{r}{2}\varphi^2X, \tag{3.11}$$

and from formula (3.7) we get

$$R(X, \xi)\xi = -\omega(X)\psi - \nabla_X\psi + e^{2\rho}\eta(X)\xi. \tag{3.12}$$

In view of (3.11) and (3.12), we obtain our formula. \square

Corollary 3.12. *In a 3-dimensional C_{12} -manifold, the Ricci tensor and the curvature tensor are given respectively by*

$$\begin{aligned} S(X, Y) &= \left(\frac{r}{2} + \operatorname{div}\psi\right)g(X, Y) + (e^{2\rho} - 2\operatorname{div}\psi - \frac{r}{2})\eta(X)\eta(Y) \\ &\quad - \omega(X)\omega(Y) - g(\nabla_X\psi, Y), \end{aligned} \tag{3.13}$$

and

$$\begin{aligned}
 R(X, Y)Z &= \left(e^{2\rho} - 2 \operatorname{div} \psi - \frac{r}{2} \right) \eta(Z)(\eta(Y)X - \eta(X)Y) \\
 &\quad - g(Y, Z) \left(\omega(X)\psi + \nabla_X \psi - \left(2 \operatorname{div} \psi + \frac{r}{2} \right) X \right) \\
 &\quad + g(X, Z) \left(\omega(Y)\psi + \nabla_Y \psi - \left(2 \operatorname{div} \psi + \frac{r}{2} \right) Y \right) \\
 &\quad + \left(e^{2\rho} - 2 \operatorname{div} \psi - \frac{r}{2} \right) \left(g(Y, Z)\eta(X) - g(X, Z)\eta(Y) \right) \xi \\
 &\quad - \omega(Z) \left(\omega(Y)X - \omega(X)Y \right) + g(\nabla_X \psi, Z)Y - g(\nabla_Y \psi, Z)X.
 \end{aligned} \tag{3.14}$$

Proof. Equation (3.13) follows from (3.9) and (3.10). Using (3.9) and (3.13) in (3.8), the curvature tensor in a 3-dimensional C_{12} -manifold is given by (3.14). \square

4. C_{12} -STRUCTURES ON THREE-DIMENSIONAL LIE GROUPS

An almost contact metric structure (φ, ξ, η, g) on a connected Lie group G is said to be left invariant if g is left invariant and if the left multiplication map $L_a : G \rightarrow G$, $L_a(x) = a.x$ has the properties

$$\varphi \circ L_a = L_a \circ \varphi \quad \text{and} \quad L_a(\xi) = \xi \quad \text{for all } a \in G.$$

Let \mathfrak{g} be an odd-dimensional Lie algebra. An almost contact metric structure on \mathfrak{g} is a quadruple (φ, ξ, η, g) , where η is a one-form, φ is an endomorphism of \mathfrak{g} and $\xi \in \mathfrak{g}$ such that

$$\eta(\xi) = 1, \quad \varphi^2(X) = -X + \eta(X)\xi, \quad g(\varphi X, \varphi Y) = g(X, Y) - \eta(X)\eta(Y)$$

for all vector fields X, Y and g is a positive definite compatible inner product on \mathfrak{g} . It is also convenient to use defining relations for the structures on Lie algebras. For instance, an almost contact metric structure (φ, ξ, η, g) on a Lie algebra \mathfrak{g} is said to be a C_{12} -structure if and only if

$$\nabla_X \xi = -\eta(X)\psi = \eta(X)\nabla_\xi \xi \tag{4.1}$$

for all X vector field in \mathfrak{g} .

Let G be a connected Lie group of dimension 3, endowed with a left invariant almost contact metric structure (φ, ξ, η, g) and let $\mathfrak{g} \cong T_e G$ be the corresponding Lie algebra of G . If $\{e_1, e_2, e_3\}$ is an orthonormal basis on \mathfrak{g} then

$$\varphi e_i = \sum_j \varphi_i^j e_j \quad \text{and} \quad \xi = a e_1 + b e_2 + c e_3,$$

where φ_i^j and a, b, c are constants such that $a^2 + b^2 + c^2 = 1$.

A classification of the Lie algebras of dimension three is found in [8], where Patera et al. list the nine classes of three-dimensional and twelve classes of four-dimensional Lie algebras. Here is the list of non-abelian three-dimensional algebras along with their defining Lie bracket equations.

| Name | Structure equations | | |
|-------------------|-----------------------------------|----------------------------------|-----------------------|
| $A_{3,1}$ | $[e_2, e_3] = e_1$ | | |
| $A_{3,2}$ | $[e_1, e_3] = e_1$ | $[e_2, e_3] = e_1 + e_2$ | |
| $A_{3,3}$ | $[e_1, e_3] = e_1$ | $[e_2, e_3] = e_2$ | |
| $A_{3,4}$ | $[e_1, e_3] = e_1$ | $[e_2, e_3] = -e_2$ | |
| $A_{3,5}^\lambda$ | $[e_1, e_3] = e_1$ | $[e_2, e_3] = \lambda e_2$ | $(0 < \lambda < 1)$ |
| $A_{3,6}$ | $[e_1, e_3] = -e_2$ | $[e_2, e_3] = e_1$ | |
| $A_{3,7}^\lambda$ | $[e_1, e_3] = -\lambda e_1 - e_2$ | $[e_2, e_3] = e_1 + \lambda e_2$ | $(\lambda > 0)$ |
| $A_{3,8}$ | $[e_1, e_2] = e_1$ | $[e_1, e_3] = -2e_2$ | $[e_2, e_3] = e_3$ |
| $A_{3,9}$ | $[e_1, e_2] = e_3$ | $[e_1, e_3] = -e_2$ | $[e_2, e_3] = e_1$ |

We will investigate the existence of C_{12} -structures on each $A_{3,i}$ and it is sufficient here to find ξ and ψ . From (4.1), we conclude that the existence of the C_{12} -structure is equivalent to

$$\nabla_{e_i} \xi = g(\xi, e_i) \nabla_\xi \xi$$

for any $i \in \{1, 2, 3\}$ or equivalently,

$$\begin{cases} \nabla_{e_1} \xi = a \nabla_\xi \xi \\ \nabla_{e_2} \xi = b \nabla_\xi \xi \\ \nabla_{e_3} \xi = c \nabla_\xi \xi. \end{cases} \tag{4.2}$$

In other words, the existence of C_{12} -structures requires the existence of the constants a, b and c provided that $\nabla_\xi \xi \neq 0$.

The algebra $A_{3,1}$. By Koszul's formula, the covariant derivatives of the basis elements are as follows:

$$\begin{aligned} \nabla_{e_1} e_1 &= 0 & \nabla_{e_1} e_2 &= -\frac{1}{2} e_3 & \nabla_{e_1} e_3 &= \frac{1}{2} e_2 \\ \nabla_{e_2} e_1 &= -\frac{1}{2} e_3 & \nabla_{e_2} e_2 &= 0 & \nabla_{e_2} e_3 &= \frac{1}{2} e_1 \\ \nabla_{e_3} e_1 &= \frac{1}{2} e_2 & \nabla_{e_3} e_2 &= -\frac{1}{2} e_1 & \nabla_{e_3} e_3 &= 0. \end{aligned}$$

By a simple computation using the covariant derivatives of the basis elements, one can get

$$\nabla_{e_1} \xi = \begin{pmatrix} 0 \\ \frac{c}{2} \\ -\frac{b}{2} \end{pmatrix}, \nabla_{e_2} \xi = \begin{pmatrix} \frac{c}{2} \\ 0 \\ -\frac{a}{2} \end{pmatrix}, \nabla_{e_3} \xi = \begin{pmatrix} -\frac{b}{2} \\ \frac{a}{2} \\ 0 \end{pmatrix} \text{ and } \nabla_\xi \xi = \begin{pmatrix} 0 \\ ac \\ -ab \end{pmatrix}.$$

With the help of system (4.2), we obtain

$$a = b = c = 0.$$

Then, there exists no C_{12} -structure on $A_{3,1}$.

The algebra $A_{3,2}$. By Koszul's formula, the covariant derivatives of the basis elements are as follows:

$$\begin{aligned} \nabla_{e_1} e_1 &= -e_3 & \nabla_{e_1} e_2 &= -\frac{1}{2}e_3 & \nabla_{e_1} e_3 &= e_1 + \frac{1}{2}e_2 \\ \nabla_{e_2} e_1 &= -\frac{1}{2}e_3 & \nabla_{e_2} e_2 &= -e_3 & \nabla_{e_2} e_3 &= \frac{1}{2}e_1 + e_2 \\ \nabla_{e_3} e_1 &= \frac{1}{2}e_2 & \nabla_{e_3} e_2 &= -\frac{1}{2}e_1 & \nabla_{e_3} e_3 &= 0. \end{aligned}$$

One can get

$$\nabla_{e_1} \xi = \begin{pmatrix} c \\ \frac{c}{2} \\ -a - \frac{b}{2} \end{pmatrix}, \quad \nabla_{e_2} \xi = \begin{pmatrix} \frac{c}{2} \\ c \\ -\frac{a}{2} - b \end{pmatrix}, \quad \nabla_{e_3} \xi = \begin{pmatrix} -\frac{b}{2} \\ \frac{a}{2} \\ 0 \end{pmatrix}$$

and

$$\nabla_{\xi} \xi = \begin{pmatrix} ac \\ ac + bc \\ -a^2 - b^2 - ab \end{pmatrix}.$$

With the help of system 4.2, we get

$$a = b = c = 0.$$

Then, there exists no C_{12} -structure on $A_{3,2}$.

The algebra $A_{3,3}$. By Koszul's formula, the covariant derivatives of the basis elements are as follows:

$$\begin{aligned} \nabla_{e_1} e_1 &= -e_3 & \nabla_{e_1} e_2 &= 0 & \nabla_{e_1} e_3 &= e_1 \\ \nabla_{e_2} e_1 &= 0 & \nabla_{e_2} e_2 &= -e_3 & \nabla_{e_2} e_3 &= e_2 \\ \nabla_{e_3} e_1 &= 0 & \nabla_{e_3} e_2 &= 0 & \nabla_{e_3} e_3 &= 0. \end{aligned}$$

One can get

$$\nabla_{e_1} \xi = \begin{pmatrix} c \\ 0 \\ -a \end{pmatrix}, \quad \nabla_{e_2} \xi = \begin{pmatrix} 0 \\ c \\ -b \end{pmatrix}, \quad \nabla_{e_3} \xi = \begin{pmatrix} 0 \\ 0 \\ 0 \end{pmatrix} \quad \text{and} \quad \nabla_{\xi} \xi = \begin{pmatrix} ac \\ bc \\ -a^2 - b^2 \end{pmatrix}.$$

With the help of system 4.2, we get an infinite number of solutions of the form

$$c = 0 \quad \text{with} \quad a^2 + b^2 = 1,$$

i.e.,

$$\xi = ae_1 \pm \sqrt{1 - a^2}e_2, \quad \text{with} \quad a \in [-1, +1] \quad \text{and} \quad \psi = e_3.$$

Then, there exists an infinite number of C_{12} -structures on $A_{3,3}$.

The algebra $A_{3,4}$. By Koszul's formula, the covariant derivatives of the basis elements are as follows:

$$\begin{aligned} \nabla_{e_1} e_1 &= -e_3 & \nabla_{e_1} e_2 &= 0 & \nabla_{e_1} e_3 &= e_1 \\ \nabla_{e_2} e_1 &= 0 & \nabla_{e_2} e_2 &= e_3 & \nabla_{e_2} e_3 &= -e_2 \\ \nabla_{e_3} e_1 &= 0 & \nabla_{e_3} e_2 &= 0 & \nabla_{e_3} e_3 &= 0. \end{aligned}$$

Therefore, we obtain

$$\nabla_{e_1}\xi = \begin{pmatrix} c \\ 0 \\ -a \end{pmatrix}, \nabla_{e_2}\xi = \begin{pmatrix} 0 \\ c \\ -b \end{pmatrix}, \nabla_{e_3}\xi = \begin{pmatrix} 0 \\ 0 \\ 0 \end{pmatrix} \text{ and } \nabla_\xi\xi = \begin{pmatrix} ac \\ bc \\ -a^2 - b^2 \end{pmatrix}.$$

With the help of system 4.2, we get four solutions of the form

$$(a, b, c) \in \{(1, 0, 0); (-1, 0, 0); (0, 1, 0); (0, -1, 0)\},$$

i.e.,

$$(\xi, \psi) \in \{(e_1, e_3), (-e_1, e_3), (e_2, e_3), (-e_2, e_3)\}.$$

So, there exists an infinite number of C_{12} -structures on $A_{3,4}$ with

$$(\xi, \psi) \in \{(e_1, e_3), (-e_1, e_3), (e_2, e_3), (-e_2, e_3)\} \text{ and } \varphi e_i = \sum_j \varphi_i^j e_j.$$

The algebra $A_{3,5}^\lambda$. By Koszul’s formula, the covariant derivatives of the basis elements are as follows:

$$\begin{aligned} \nabla_{e_1}e_1 &= -e_3 & \nabla_{e_1}e_2 &= 0 & \nabla_{e_1}e_3 &= e_1 \\ \nabla_{e_2}e_1 &= 0 & \nabla_{e_2}e_2 &= -\lambda e_3 & \nabla_{e_2}e_3 &= \lambda e_2 \\ \nabla_{e_3}e_1 &= 0 & \nabla_{e_3}e_2 &= 0 & \nabla_{e_3}e_3 &= 0. \end{aligned}$$

Therefore, we obtain

$$\nabla_{e_1}\xi = \begin{pmatrix} c \\ 0 \\ -a \end{pmatrix}, \nabla_{e_2}\xi = \begin{pmatrix} 0 \\ \lambda c \\ -\lambda b \end{pmatrix}, \nabla_{e_3}\xi = \begin{pmatrix} 0 \\ 0 \\ 0 \end{pmatrix} \text{ and } \nabla_\xi\xi = \begin{pmatrix} ac \\ \lambda bc \\ -a^2 - \lambda b^2 \end{pmatrix}.$$

Replacing in the system 4.2, we get four solutions of the form

$$(a, b, c) \in \{(1, 0, 0); (-1, 0, 0); (0, 1, 0); (0, -1, 0)\},$$

i.e.,

$$(\xi, \psi) \in \{(e_1, e_3), (-e_1, e_3), (e_2, \lambda e_3), (-e_2, \lambda e_3)\}.$$

Then, there exists an infinite number of C_{12} -structures on $A_{3,5}^\lambda$ with $0 < \lambda < 1$.

The algebra $A_{3,6}$. By Koszul’s formula, the covariant derivatives of the basis elements are as follows:

$$\begin{aligned} \nabla_{e_1}e_1 &= 0 & \nabla_{e_1}e_2 &= 0 & \nabla_{e_1}e_3 &= 0 \\ \nabla_{e_2}e_1 &= 0 & \nabla_{e_2}e_2 &= 0 & \nabla_{e_2}e_3 &= 0 \\ \nabla_{e_3}e_1 &= e_2 & \nabla_{e_3}e_2 &= -e_1 & \nabla_{e_3}e_3 &= 0. \end{aligned}$$

One can get

$$\nabla_{e_1}\xi = \nabla_{e_2}\xi = \begin{pmatrix} 0 \\ 0 \\ 0 \end{pmatrix}, \nabla_{e_3}\xi = \begin{pmatrix} -b \\ a \\ 0 \end{pmatrix} \text{ and } \nabla_\xi\xi = \begin{pmatrix} -bc \\ ac \\ 0 \end{pmatrix}.$$

From system 4.2, we get $a = b = 0$ and $c \in \mathbb{R}$ this implies $\nabla_\xi\xi = 0$. Then, there exists no C_{12} -structure on $A_{3,6}$.

The algebra $A_{3,7}^\lambda$. By Koszul's formula, the covariant derivatives of the basis elements are as follows:

$$\begin{aligned} \nabla_{e_1} e_1 &= \lambda e_3 & \nabla_{e_1} e_2 &= 0 & \nabla_{e_1} e_3 &= -\lambda e_1 \\ \nabla_{e_2} e_1 &= 0 & \nabla_{e_2} e_2 &= -\lambda e_3 & \nabla_{e_2} e_3 &= \lambda e_2 \\ \nabla_{e_3} e_1 &= e_2 & \nabla_{e_3} e_2 &= -e_1 & \nabla_{e_3} e_3 &= 0. \end{aligned}$$

One can get

$$\nabla_{e_1} \xi = \lambda \begin{pmatrix} -c \\ 0 \\ a \end{pmatrix}, \quad \nabla_{e_2} \xi = \lambda \begin{pmatrix} 0 \\ c \\ -b \end{pmatrix}, \quad \nabla_{e_3} \xi = \begin{pmatrix} -b \\ a \\ 0 \end{pmatrix}$$

and

$$\nabla_\xi \xi = \begin{pmatrix} -c(a\lambda + b) \\ c(a + b\lambda) \\ \lambda(a^2 - b^2) \end{pmatrix}.$$

From 4.2, we get

$$a = b = c = 0.$$

Then, there exists no C_{12} -structure on $A_{3,7}^\lambda$.

The algebra $A_{3,8}$. By Koszul's formula, the covariant derivatives of the basis elements are as follows:

$$\begin{aligned} \nabla_{e_1} e_1 &= -e_2 & \nabla_{e_1} e_2 &= e_1 + e_3 & \nabla_{e_1} e_3 &= -e_2 \\ \nabla_{e_2} e_1 &= e_3 & \nabla_{e_2} e_2 &= 0 & \nabla_{e_2} e_3 &= -e_1 \\ \nabla_{e_3} e_1 &= e_2 & \nabla_{e_3} e_2 &= -e_1 - e_3 & \nabla_{e_3} e_3 &= e_2. \end{aligned}$$

One can get

$$\nabla_{e_1} \xi = -\nabla_{e_3} \xi = \begin{pmatrix} b \\ -a - c \\ b \end{pmatrix}, \quad \nabla_{e_2} \xi = \begin{pmatrix} -c \\ 0 \\ a \end{pmatrix} \quad \text{and} \quad \nabla_\xi \xi = \begin{pmatrix} b(a - 2c) \\ -a^2 + c^2 \\ b(2a - c) \end{pmatrix}.$$

From 4.2, we obtain the system

$$a^2 = b^2 = \frac{1}{3} \quad \text{and} \quad c = -a,$$

which gives four solutions;

$$(a, b, c) \in \left\{ \frac{1}{\sqrt{3}}(1, 1, -1); \frac{1}{\sqrt{3}}(1, -1, -1); \frac{1}{\sqrt{3}}(-1, 1, 1); \frac{1}{\sqrt{3}}(-1, -1, 1) \right\}.$$

So, there exists an infinite number of C_{12} -structures on $A_{3,8}$.

The algebra $A_{3,9}$. By Koszul's formula, the covariant derivatives of the basis elements are as follows:

$$\begin{aligned}\nabla_{e_1}e_1 &= 0 & \nabla_{e_1}e_2 &= \frac{1}{2}e_3 & \nabla_{e_1}e_3 &= -\frac{1}{2}e_2 \\ \nabla_{e_2}e_1 &= -\frac{1}{2}e_3 & \nabla_{e_2}e_2 &= 0 & \nabla_{e_2}e_3 &= \frac{1}{2}e_1 \\ \nabla_{e_3}e_1 &= \frac{1}{2}e_2 & \nabla_{e_3}e_2 &= \frac{1}{2}e_1 & \nabla_{e_3}e_3 &= 0.\end{aligned}$$

By a simple computation using the covariant derivatives of the basis elements, one can get

$$\nabla_{e_1}\xi = \begin{pmatrix} 0 \\ -\frac{c}{2} \\ \frac{b}{2} \end{pmatrix}, \nabla_{e_2}\xi = \begin{pmatrix} \frac{c}{2} \\ 0 \\ -\frac{a}{2} \end{pmatrix}, \nabla_{e_3}\xi = \begin{pmatrix} -\frac{b}{2} \\ \frac{a}{2} \\ 0 \end{pmatrix} \text{ and } \nabla_{\xi}\xi = \begin{pmatrix} 0 \\ 0 \\ 0 \end{pmatrix}.$$

Since $\nabla_{\xi}\xi = 0$, there exists no C_{12} -structure on $A_{3,9}$.

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Gherici Beldjilali

Laboratory of Quantum Physics and Mathematical Modeling (LPQ3M), University of Mascara, Algeria
gherici.beldjilali@univ-mascara.dz

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