

SOBOLEV SPACES DIVERSIFICATION

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ABSTRACT. This work attempts to be an overview of a variety of results concerning Sobolev spaces associated to some orthonormal systems, particularly the Hermite and Laguerre operators settings.

1. INTRODUCTION

This work is based on a talk given in “X Encuentro de Analistas A. P. Calderón” at La Falda, Córdoba, Argentina, August 25-28, 2010. The talk was about some results obtained in [5] and [6] about Sobolev spaces associated to the Hermite and Laguerre operators.

Hermite-Sobolev spaces already appear in [23] for the case $p = 2$ where an expansion type definition was used. For $p > 1$ an approach considering *derivatives* was presented in [5]. These results were extended in [6] and the new ideas were used to describe Laguerre-Sobolev spaces.

Sobolev spaces associated to other operators were also studied in the last years creating an interesting diversity. Examples of these are [2], [3], [11] and [15] among others.

In the classical theory, given a multi-index $\alpha = (\alpha_j)_{j=1}^d$ of non-negative integers, we denote the operator

$$\frac{\partial^\alpha}{\partial x^\alpha} = \frac{\partial^{|\alpha|}}{\partial x_1^{\alpha_1} \partial x_2^{\alpha_2} \dots \partial x_d^{\alpha_d}},$$

where $|\alpha| = \alpha_1 + \alpha_2 + \dots + \alpha_n$, and the derivatives $\frac{\partial}{\partial x_i}$, $i = 1, \dots, d$, are taken in the weak sense. Then, for $1 \leq p \leq \infty$, the classical *Sobolev space* of order $k \in \mathbb{N}_0$ is defined by

$$W^{k,p} = \{f \in L^p(\mathbb{R}^d) : \frac{\partial^\alpha}{\partial x^\alpha} f \in L^p(\mathbb{R}^d), |\alpha| \leq k\}.$$

The space $W^{k,p}$ is a separable Banach space with the norm

$$\|f\|_{k,p} = \sum_{|\alpha| \leq k} \left\| \frac{\partial^\alpha}{\partial x^\alpha} f \right\|_p,$$

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where $\|\cdot\|_p$ denotes the usual $L^p(\mathbb{R}^d)$ norm. It is also well known that the set C_c^∞ (the set of functions with infinitely many derivatives and compact support) is a dense subspace (see [18]).

If we start with $f \in L^2(\mathbb{R}^d)$, by the formula

$$\left(\frac{\partial}{\partial \zeta_j} f\right)^\wedge(\zeta) = -2\pi i \zeta_j \widehat{f}(\zeta),$$

we have $f \in W^{k,2}$ if and only if $(1 + |\zeta|^2)^{k/2} \widehat{f} \in L^2(\mathbb{R}^d)$.

Therefore, it is reasonable to define for $s > 0$

$$W^{s,p} = \left\{ f \in L^2(\mathbb{R}^d) : \int_{\mathbb{R}^d} |\widehat{f}(\zeta)|^2 (1 + |\zeta|^2)^s d\zeta < \infty \right\}.$$

Now, using

$$[(I - \Delta)f]^\wedge(\zeta) = (1 + 4\pi^2|\zeta|^2)\widehat{f},$$

we have for $s > 0$,

$$[(I - \Delta)^{-s/2} f]^\wedge(\zeta) = (1 + 4\pi^2|\zeta|^2)^{-s/2} \widehat{f}(\zeta).$$

The operator $(I - \Delta)^{-s/2}$, with $s > 0$, is called the Bessel potential of order s . Therefore, $f \in W^{s,2}$ if and only if there exists $g \in L^2(\mathbb{R}^d)$ such that $(I - \Delta)^{-s/2} g = f$, that is to say

$$W^{2,p} = (I - \Delta)^{-s/2}(L^2(\mathbb{R}^d)).$$

Bessel potentials are bounded operators on $L^p(\mathbb{R}^d)$, $p \geq 1$ (see [18], Chapter V).

For $p \geq 1$ and $s > 0$, the *potential space* of order s and integrability p is defined as

$$L_s^p = (I - \Delta)^{-s/2}(L^p(\mathbb{R}^d))$$

with the norm

$$\|f\|_{L_s^p} = \|g\|_p,$$

where g is such that $(I - \Delta)^{-s/2} g = f$.

In [18] it was proven that for a positive integer k and $1 < p < \infty$, the space $W^{k,p} = L_k^p$.

A more general setting might be constructed considering a second order differential operator \mathbf{L} self-adjoint with respect to a measure μ .

Sometimes, it is possible to obtain a factorization of \mathbf{L} as

$$\mathbf{L} = \sum_i \partial_i \tilde{\partial}_i,$$

where ∂_i and $\tilde{\partial}_i$ are first order differential operators, and then to define the *Sobolev space* of order k and integrability p associated to \mathbf{L} as

$$W^{k,p} = \{f \in L^p(\mu) : (\partial_i)^j f \in L^p(\mu) \text{ and } (\tilde{\partial}_i)^j f \in L^p(\mu), j \leq k\}.$$

When $s > 0$ the definition of the *potential space* of order s and integrability p could be given by

$$L_s^p = (I + \mathbf{L})^{-s/2}(L^p(\mu)).$$

In the case that $\mathbf{L}^{-s/2}$ is bounded on $L^p(\mu)$, we have

$$L_s^p = \mathbf{L}^{-s/2}(L^p(\mu)). \tag{1}$$

It is reasonable to expect that for a positive integer k it follows $W^{k,p} = L_k^p$, and this tell us that ∂_i and $\tilde{\partial}_i$ have the right to be called *derivatives* associated to \mathbf{L} .

2. HERMITE SETTING

The Hermite operator is defined as

$$H = -\Delta + |x|^2, \quad x \in \mathbb{R}^d. \tag{2}$$

Its eigenvectors are the Hermite functions, which form an orthonormal basis for $L^2(\mathbb{R}^d)$. In \mathbb{R} they are defined for $n \in \mathbb{N}_0 = \{0, 1, 2, 3, \dots\}$ as

$$h_n(t) = \frac{H_n(t) e^{-t^2/2}}{(2^n n! \pi^{1/2})^{1/2}}, \quad t \in \mathbb{R},$$

where H_n is the Hermite polynomial of order n (see [21]).

In \mathbb{R}^d given a multi-index $\alpha = (\alpha_j)_{j=1}^d \in \mathbb{N}_0^d$, the Hermite function of order α is defined using the unidimensional ones by

$$h_\alpha(x) = \prod_{j=1}^d h_{\alpha_j}(x_j), \quad x = (x_1, \dots, x_d) \in \mathbb{R}^d.$$

As we said, the Hermite functions are eigenvectors of H (see [22]) satisfying

$$Hh_\alpha = (2|\alpha| + d) h_\alpha,$$

where $|\alpha| = \sum_{j=1}^d \alpha_j$.

Some other interesting properties of Hermite functions are the following. Their proofs can be found in [19] and [22].

Proposition 1. *If $M \in \mathbb{N}$ and $f \in C_c^\infty$, then there exists a constant $C_{M,f} > 0$ such that*

$$\left| \int_{\mathbb{R}^d} f h_\alpha \right| \leq C_{M,f} (|\alpha| + 1)^{-M}, \quad \alpha \in \mathbb{N}^d.$$

Proposition 2. *If $1 \leq p < \infty$ and $w \in A_p$, then there exist constants $\epsilon_p > 0$ and C_w such that*

$$\|h_\alpha\|_{L^p(w)} \leq C_w (|\alpha| + 1)^{\epsilon_p}.$$

Proposition 3. *As $n \rightarrow \infty$ the Hermite functions satisfy the estimates:*

- i) $\|h_n\|_p \sim n^{\frac{1}{2p} - \frac{1}{4}}, \quad 1 \leq p < 4,$
- ii) $\|h_n\|_p \sim n^{-\frac{1}{8}} \log(n), \quad p = 4,$
- iii) $\|h_n\|_p \sim n^{\frac{1}{6p} - \frac{1}{12}}, \quad 4 < p \leq \infty.$

The Hermite operator can be factorized as

$$H = \frac{1}{2} \sum_{j=1}^d A_j A_{-j} + A_{-j} A_j, \tag{3}$$

where

$$A_j = \frac{\partial}{\partial x_j} + x_j \quad \text{and} \quad A_{-j} = A_j^* = -\frac{\partial}{\partial x_j} + x_j, \quad 1 \leq j \leq d.$$

The operators A_j and A_{-j} , are called *annihilation* and *creation* operators respectively because of their behavior over Hermite functions, in fact, for $1 \leq j \leq d$,

$$A_j h_\alpha = \sqrt{2\alpha_j} h_{\alpha - e_j}, \quad A_{-j} h_\alpha = \sqrt{2(\alpha_j + 1)} h_{\alpha + e_j}, \quad (4)$$

where e_j is the j th-coordinate vector in \mathbb{N}_0^d . In the context of the Hermite operator, the notion of *derivatives* is given by these operators.

Given $p \in (1, \infty)$ and $k \in \mathbb{N}$, the Hermite-Sobolev space of order k , denoted by $W^{k,p}$, is the set of functions $f \in L^p(\mathbb{R}^d)$ such that

$$A_{j_1} \cdots A_{j_m} f \in L^p(\mathbb{R}^d), \quad 1 \leq |j_1|, \dots, |j_m| \leq d, \quad 1 \leq m \leq k,$$

with the norm

$$\|f\|_{W^{k,p}} = \sum_{1 \leq |j_1|, \dots, |j_m| \leq d, 1 \leq m \leq k} \|A_{j_1} \cdots A_{j_m} f\|_p + \|f\|_p.$$

This definition was rewritten in [6] considering only annihilation operators proving they are enough to define the Hermite-Sobolev space. In the same work, the authors deal with weighted Sobolev spaces for weights in the Muckenhoupt class A_p , defined for $1 < p < \infty$, as the set of weights (non-negative and locally integrable functions) w such that

$$\left(\int_B w \right)^{1/p} \left(\int_B w^{-\frac{1}{p-1}} \right)^{1/p'} \leq C|B|, \quad (5)$$

for every ball $B \subset \mathbb{R}^d$; and for $p = 1$, A_1 is defined as those weights w satisfying the condition

$$w(B) \sup_B w^{-1} \leq C|B|, \quad (6)$$

for every ball $B \subset \mathbb{R}^d$.

Given a weight w , $k \in \mathbb{N}$ and $p \geq 1$, the Hermite-Sobolev space of order k , denoted by $W^{k,p}(w)$, is defined as the set of functions $f \in L^p(w)$ such that

$$\overbrace{A_j \cdots A_j}^{m \text{ times}} f = A_j^m f \in L^p(w), \quad 1 \leq m \leq k, \quad 1 \leq j \leq d,$$

with the norm

$$\|f\|_{W^{k,p}(w)} = \sum_{j=1}^d \sum_{1 \leq m \leq k} \|A_j^m f\|_{L^p(w)} + \|f\|_{L^p(w)}.$$

On the other hand, in order to define a potential space like (1), for $a > 0$ we define the operator

$$H^{-a} f(x) = \frac{1}{\Gamma(a)} \int_0^\infty e^{-tH} f(x) t^a \frac{dt}{t}, \quad x \in \mathbb{R}^d, \quad f \in \mathfrak{F}, \quad (7)$$

where $\{e^{-tH}\}_{t \geq 0}$ is the heat semi-group associated to H , and \mathfrak{F} denotes the set of linear combinations of Hermite functions.

Remark 1. If $a > 0$ and $\alpha \in \mathbb{N}_0^d$, by using the Γ function and the fact

$$e^{-tH} h_\alpha = e^{-t(2|\alpha|+d)} h_\alpha,$$

we have

$$H^{-a} h_\alpha(x) = \frac{1}{\Gamma(a)} \int_0^\infty e^{-tH} h_\alpha(x) t^a \frac{dt}{t} = (2|\alpha| + d)^{-a} h_\alpha(x), \quad x \in \mathbb{R}^d.$$

The operator H^{-a} has a kernel K_a with exponential behavior far from the diagonal as the following proposition shows. The proof can be found in [5].

Proposition 4. *The operator H^{-s} , $s > 0$, has an integral representation*

$$H^{-s} f(x) = \int_{\mathbb{R}^d} K_s(x, y) f(y) dy, \quad x \in \mathbb{R}^d,$$

where $K_s(x, y)$ is positive and symmetric. Moreover,

$$K_a(x, y) \leq C \phi_a(|x - y|), \quad x, y \in \mathbb{R}^d, \tag{8}$$

where $\phi_a(r)$, for $r \geq 0$, is defined by

$$\phi_a(r) = \begin{cases} \frac{\chi_{\{r < 1\}}(r)}{r^{d-2a}} + e^{-\frac{r^2}{4}} \chi_{\{r \geq 1\}}(r), & \text{if } a < \frac{d}{2}, \\ \log\left(\frac{e}{r}\right) \chi_{\{r < 1\}}(r) + e^{-\frac{r^2}{4}} \chi_{\{r \geq 1\}}(r), & \text{if } a = \frac{d}{2}, \\ \chi_{\{r < 1\}}(r) + e^{-\frac{r^2}{4}} \chi_{\{r \geq 1\}}(r), & \text{if } a > \frac{d}{2}. \end{cases}$$

In [6] it was proven that H^{-a} a bounded operator on $L^p(w)$ for a weight $w \in A_p$ as the following theorem states.

Theorem 1. *Let $1 \leq p < \infty$ and $a > 0$. If $w \in A_p$, then the operator H^{-a} is bounded on $L^p(w)$.*

Remark 2. The results in [4] (Theorem 4 therein) suggest that there should be more weights for the boundedness of H^{-a} in $L^p(w)$. Those classes of weights allow power weights $w(x) = |x|^\gamma$ without restriction on γ .

For the unweighted case L^p - L^q inequalities were obtained in [5].

Theorem 2. *Let a, d such that $0 < a < d$, then:*

i) *There exists a constant C such that*

$$\|H^{-a/2} f\|_q \leq C \|f\|_1,$$

for all $f \in L^1(\mathbb{R}^d)$ if and only if $1 \leq q < \frac{d}{d-a}$.

ii) *There exists a constant C such that*

$$\|H^{-a/2} f\|_\infty \leq C \|f\|_p,$$

for all f in $L^p(\mathbb{R}^d)$ if and only if $p > \frac{d}{a}$.

iii) There exists a constant C such that

$$\|H^{-a/2}f\|_q \leq C\|f\|_\infty,$$

for all $f \in L^\infty(\mathbb{R}^d)$ if and only if $q > \frac{d}{a}$.

iv) There exists a constant C such that

$$\|H^{-a/2}f\|_1 \leq C\|f\|_p,$$

for all $f \in L^p(\mathbb{R}^d)$ if and only if $1 \leq p < \frac{d}{d-a}$.

v) If $1 < p < \infty$, $1 < q < \infty$ and $\frac{1}{p} - \frac{a}{d} \leq \frac{1}{q} < \frac{1}{p} + \frac{a}{d}$, then there exists a constant C such that

$$\|H^{-a/2}f\|_q \leq C\|f\|_p,$$

for all $f \in L^p(\mathbb{R}^d)$.

In the Hermite setting the potential space of order a and integrability p is defined as $\mathcal{L}_a^p(w) = H^{-a/2}(L^p(w))$, with respect to an absolute continuous measure $w(x)dx$, being w a weight in a class where $H^{-a/2}$ results bounded. From Theorem 1 it is enough to ask $w \in A_p$.

A norm on $\mathcal{L}_a^p(w)$ is given by

$$\|f\|_{\mathcal{L}_a^p(w)} = \|g\|_{L^p(w)},$$

where g is such that $H^{-a/2}g = f$.

Remark 3. It is easy to see that $H^{-a/2}$ is one to one (see [5]), and this assures that the space \mathfrak{L}_a^p is well defined for $p \in [1, \infty)$ and $a > 0$. As $\mathfrak{F} = H^{-a/2}(\mathfrak{F})$ then \mathfrak{F} is a dense space of \mathfrak{L}_a^p .

A fundamental devise of the theory is the family of Hermite-Riesz transforms of order $m \in \mathbb{N}$, associated to H , defined by

$$R_J^m = A_{j_1} \dots A_{j_m} H^{-m/2}, \text{ where } J = (j_1, \dots, j_m), 1 \leq |j_i| \leq d, 1 \leq i \leq m.$$

In the case $j_1 = \dots = j_m = j$, these operators will be denoted by R_j^m . The case $m = 1$ was introduced by S. Thangavelu (see [22]). He proved that they are bounded operators in $L^p(\mathbb{R}^d)$. Also in [19] and [20], it was shown that the operators R_j^m are Calderón-Zygmund operators and as a consequence they are bounded in $L^p(w)$ for $w \in A_p$, $1 < p < \infty$.

We shall now present some expected properties of the spaces $\mathfrak{L}_a^p(w)$ appearing in [6].

Theorem 3. *Let $w \in A_p$, $1 < p < \infty$, and $a > 0$.*

- i) *If $t > a$, then $\mathfrak{L}_t^p(w) \subset \mathfrak{L}_a^p(w) \subset L^p(w)$ with continuous inclusions. Moreover, $\mathfrak{L}_a^p(w)$ and $\mathfrak{L}_t^p(w)$ are isometrically isomorphic.*
- ii) *If $t > 0$, then $H^{-t/2}$ maps $\mathfrak{L}_a^p(w)$ into $\mathfrak{L}_{a+t}^p(w)$.*
- iii) *If $a > 1$ and $1 \leq |j| \leq d$, then A_j is bounded from $\mathfrak{L}_a^p(w)$ into $\mathfrak{L}_{a-1}^p(w)$.*
- iv) *The operators R_j^m are bounded on $\mathfrak{L}_a^p(w)$.*

For the unweighted case some comparison with the classical Sobolev spaces is presented in the following result.

Theorem 4. *Let $a > 0$ and $p \in (1, \infty)$. Then*

- i) $\mathfrak{L}_a^p \subset L_a^p$.*
- ii) $\mathfrak{L}_a^p \neq L_a^p$.*
- iii) If $f \in L_a^p$ and has compact support, then f belongs to \mathfrak{L}_a^p .*

Remark 4. The results in [10] about Sobolev spaces associated to Schrödinger operators with a polynomial potential, implies in particular that f belongs to \mathfrak{L}_a^p if and only if $f \in L_a^p$ and $|x|^{2a} f \in L^p$.

The following structural theorem was proved in [6]. A fundamental part of the proof is due to the boundedness of higher order Riesz transforms given in Theorem 3.

Theorem 5. *Let $k \in \mathbb{N}$, $1 < p < \infty$, and $w \in A_p$. Then,*

$$W^{k,p}(w) = \mathfrak{L}_k^p(w)$$

and the norms $\|\cdot\|_{W^{k,p}(w)}$ and $\|\cdot\|_{\mathfrak{L}_k^p(w)}$ are equivalent.

3. LAGUERRE SETTING

For $\alpha > -1$ and $n \in \mathbb{N}_0$, the Laguerre polynomial of order n and type α , is defined by

$$L_n^\alpha = \frac{x^\alpha e^x}{n!} \frac{d^n}{dx^n} (e^{-x} x^{n+\alpha}),$$

and the Laguerre function of order n and type α is

$$\mathcal{L}_n^\alpha(y) = \left(\frac{\Gamma(n+1)}{\Gamma(n+\alpha+1)} \right)^{1/2} e^{-y/2} y^{\alpha/2} L_n^\alpha(y), \quad y \in \mathbb{R}^+, \quad n \in \mathbb{N}_0. \quad (9)$$

For each $\alpha > -1$, $\{\mathcal{L}_n^\alpha\}_{n=0}^\infty$ is an orthonormal system in $L^2((0, \infty))$ and satisfies

$$L_\alpha \mathcal{L}_n^\alpha = \left(n + \frac{\alpha+1}{2} \right) \mathcal{L}_n^\alpha, \quad n \in \mathbb{N}_0,$$

where L_α is the Laguerre operator, self-adjoint in the set $\mathcal{C}_c(0, \infty)$, defined by

$$L_\alpha = -y \frac{d^2}{dy^2} - \frac{d}{dy} + \frac{y}{4} + \frac{\alpha^2}{4y}, \quad y \in (0, \infty). \quad (10)$$

The operator L_α can be written as

$$L_\alpha = (\delta^\alpha)^* \delta^\alpha + \frac{(\alpha+1)}{2},$$

where

$$\delta^\alpha = \sqrt{x} \frac{d}{dx} + \frac{1}{2} \left(\sqrt{x} - \frac{\alpha}{\sqrt{x}} \right) \quad \text{and} \quad (\delta^\alpha)^* = -\sqrt{x} \frac{d}{dx} + \frac{1}{2} \left(\sqrt{x} - \frac{\alpha+1}{\sqrt{x}} \right). \quad (11)$$

As was shown in [6] the operator δ^α plays the role of *derivative* in the Laguerre setting. The action of these operators on Laguerre functions is given by

$$\delta^\alpha(\mathcal{L}_0^\alpha) = 0, \quad \delta^\alpha(\mathcal{L}_n^\alpha) = -\sqrt{n} \mathcal{L}_{n-1}^{\alpha+1}, \quad \text{for } n \geq 1, \quad \text{and} \quad (12)$$

$$(\delta^\alpha)^*(\mathcal{L}_n^{\alpha+1}) = -\sqrt{n+1} \mathcal{L}_{n+1}^\alpha \quad \text{for } n \geq 0. \quad (13)$$

In this section we will deal with power weights y^γ in order to proof boundedness on $L^p(\mathbb{R}^+, y^\gamma dy)$ of some operators associated to L_α , (see [1]). The range for the exponent γ will be

$$-\frac{\alpha}{2}p - 1 < \gamma < p - 1 + \frac{\alpha}{2}p, \tag{14}$$

where $\alpha > -1, 1 < p < \infty$.

Under this hypothesis it is known (see [21] Theorem 5.7.1) that the set S_α of finite linear combinations of Laguerre functions is dense in $L^p((0, \infty), y^\gamma dy)$.

At first sight, the natural way of defining a Sobolev space should be, iterating the derivative δ^α , as the set of functions f in $L^p(\mathbb{R}^+, y^\gamma dy)$ such that $(\delta^\alpha)^m f \in L^p(\mathbb{R}^+, y^\gamma dy), 0 \leq m \leq k$. We shall denote this space by $\mathcal{W}_\alpha^{k,p}(y^\gamma)$, and the norm is

$$\|f\|_{\mathcal{W}_\alpha^{k,p}(y^\gamma)} = \sum_{m=0}^k \|(\delta^\alpha)^m f\|_{L^p(\mathbb{R}^+, y^\gamma dy)}.$$

We will see later that this space is not appropriate as a Sobolev space for L_α when $-1 < \alpha < 0$.

It was found in [6] that the right space that plays the role of a Sobolev space defined via derivatives in the Laguerre setting is called Laguerre-Sobolev spaces, denoted by $W_\alpha^{k,p}(y^\gamma)$, defined by the sets of functions f in $L^p(\mathbb{R}^+, y^\gamma dy)$ such that

$$\delta^{\alpha+m} \circ \dots \circ \delta^{\alpha+1} \circ \delta^\alpha f \in L^p(\mathbb{R}^+, y^\gamma dy), 0 \leq m \leq k - 1.$$

The norm on $W_\alpha^{k,p}(y^\gamma)$ is given by

$$\|f\|_{W_\alpha^{k,p}(y^\gamma)} = \|f\|_{p,\gamma} + \sum_{m=0}^{k-1} \|\delta^{\alpha+m} \circ \dots \circ \delta^{\alpha+1} \circ \delta^\alpha f\|_{p,\gamma}.$$

These spaces are the right spaces for the theory as the following theorem shows (see [6]).

Theorem 6. *Let $\alpha > -1, 1 < p < \infty, k \in \mathbb{N}$ and γ satisfying (14). Then,*

$$W_\alpha^{k,p}(y^\gamma) = \mathfrak{W}_\alpha^{k,p}(y^\gamma),$$

and the norms are equivalent.

The Riesz transforms were defined in [12], for $\alpha > -1$, by

$$R_\alpha = \delta^\alpha (L_\alpha)^{-1/2} \quad \text{and} \quad \tilde{R}_\alpha = (\delta^\alpha)^* (L_{\alpha+1})^{-1/2}.$$

In [13] it was proved that those operators are bounded on $L^p(\mathbb{R}^+, y^\gamma dy)$ for γ satisfying (14). Given a positive integer k and $\alpha > -1$ we define the higher order Riesz transform of order k as

$$R_\alpha^k = \left(\delta^{\alpha+k-1} \circ \dots \circ \delta^{\alpha+1} \circ \delta^\alpha \right) (L_\alpha)^{-k/2}$$

and

$$\tilde{R}_\alpha^k = \left((\delta^\alpha)^* \circ (\delta^{\alpha+1})^* \circ \dots \circ (\delta^{\alpha+k-1})^* \right) (L_{\alpha+k})^{-k/2}.$$

Observe that $R_\alpha^1 = R_\alpha$ and $\tilde{R}_\alpha^1 = \tilde{R}_\alpha$.

It was proved in [6] the following boundedness result for the operators R_α^k and \tilde{R}_α^k .

Theorem 7. Let $k \in \mathbb{N}$, $1 < p < \infty$, $\alpha > -1$ and γ satisfying (14). The operators R_α^k and \tilde{R}_α^k are bounded on $L^p(\mathbb{R}^+, y^\gamma dy)$.

As in the Hermite setting we have the following structural theorem for the spaces $\mathfrak{W}_{\alpha,a}^p(y^\gamma)$.

Theorem 8. Let $\alpha > -1$, $1 < p < \infty$, $a > 0$, and γ satisfying (14).

- i) If $t > a$, then $\mathfrak{W}_{\alpha,a}^p(y^\gamma) \subset \mathfrak{W}_{\alpha,t}^p(y^\gamma) \subset L^p(y^\gamma)$ with continuous inclusions. Moreover, $\mathfrak{W}_{\alpha,a}^p(y^\gamma)$ and $\mathfrak{W}_{\alpha,t}^p(y^\gamma)$ are isometrically isomorphic .
- ii) If $t > 0$, then $(L_\alpha)^{-t/2}$ maps $\mathfrak{W}_{\alpha,a}^p(y^\gamma)$ into $\mathfrak{W}_{\alpha,a+t}^p(y^\gamma)$.
- iii) If $a > 1$, then δ^α is bounded from $\mathfrak{W}_{\alpha,a}^p(y^\gamma)$ into $\mathfrak{W}_{\alpha+1,a-1}^p(y^\gamma)$.
- iv) The operators $R_{\alpha,s}^k$ are bounded from $\mathfrak{W}_{\alpha,s}^p(y^\gamma)$ into $\mathfrak{W}_{\alpha+k,s}^p(y^\gamma)$.

Proposition 5. Let $\alpha > -1$, $1 < p < \infty$, $k \in \mathbb{N}$ and γ satisfying (14). Then S_α is a dense subspace of $W_\alpha^{k,p}(y^\gamma)$.

A Riesz transform of higher order k associated to L_α could be defined iterating the derivative δ^α as $(\delta^\alpha)^k(L_\alpha)^{-k/2}$. This operator has the same boundedness properties of R_α^k .

Theorem 9. Let $\alpha > -1$ and k a positive integer. Then the Riesz transforms $(\delta^\alpha)^k(L_\alpha)^{-k/2}$ are bounded in $L^p(y^\gamma dy)$ for γ and p satisfying (14).

When $-1 < \alpha \leq 0$, despite the boundedness of $(\delta^\alpha)^k(L_\alpha)^{-k/2}$ in $L^p(y^\gamma dy)$ for γ satisfying (14), the spaces $\mathcal{W}_{\alpha,k}^{k,p}(y^\delta)$ are different from the potential spaces $\mathfrak{W}_{\alpha,k}^p(y^\delta)$ as the following theorem shows.

Theorem 10. Let p be in the range $1 < p < \infty$.

- i) If $\alpha > -1$, and γ satisfies (14), then $\mathfrak{W}_{\alpha}^{k,p}(y^\gamma) \subset \mathcal{W}_{\alpha}^{k,p}(y^\gamma)$.
- ii) If $-1 < \alpha \leq 0$, then $\mathfrak{W}_{\alpha}^{2,2} \neq \mathcal{W}_{\alpha}^{2,2}$.
- iii) If $\alpha > 0$, and γ satisfies

$$-\frac{\alpha - 1}{2}p - 1 < \gamma < p - 1 + \frac{\alpha - 1}{2}p. \tag{15}$$

then $\mathfrak{W}_{\alpha}^{2,p}(y^\gamma) = \mathcal{W}_{\alpha}^{2,p}(y^\gamma)$,

4. OTHER LAGUERRE FUNCTION SYSTEMS

It is possible to translate the concepts and results of the previous section to other Laguerre systems. For instance, we shall consider the orthonormal system in $L^2((0, \infty), dy)$ given by $\varphi_k^\alpha(y) = \mathfrak{L}_k^\alpha(y^2)(2y)^{1/2}$, where \mathfrak{L}_k^α are the functions defined in (9). The functions φ_k^α are eigenfunctions of the operator

$$\mathbf{L}_\alpha = \frac{1}{4} \left\{ -\frac{d^2}{dy^2} + y^2 + \frac{1}{y^2} \left(\alpha^2 - \frac{1}{4} \right) \right\},$$

since

$$\mathbf{L}_\alpha(\varphi_k^\alpha) = \left(k + \frac{\alpha + 1}{2} \right) \varphi_k^\alpha.$$

The operator \mathbf{L}_α can be written as

$$\mathbf{L}_\alpha = (\mathbf{D}_\alpha)^* \mathbf{D}_\alpha - \left(\frac{\alpha + 1}{2}\right),$$

with $\mathbf{D}^\alpha = \frac{1}{2} \left\{ \frac{d}{dy} + y - \frac{1}{y} \left(\alpha + \frac{1}{2}\right) \right\}$ and $(\mathbf{D}^\alpha)^* = \frac{1}{2} \left\{ -\frac{d}{dy} + y - \frac{1}{y} \left(\alpha + \frac{1}{2}\right) \right\}$.

The operator $(\mathbf{D}^\alpha)^*$ is the formal adjoint of \mathbf{D}^α with respect to the Lebesgue measure. The behavior of those operators over φ_k^α is

$$\mathbf{D}^\alpha(\varphi_k^\alpha) = -\sqrt{k}\varphi_{k-1}^{\alpha+1} \quad \text{and} \quad (\mathbf{D}^{\beta-1})^*(\varphi_k^\beta) = -\sqrt{k+1}\varphi_{k+1}^{\beta-1}.$$

As in Section 3 the Riesz transforms can be defined as

$$\mathbf{R}_\alpha^k = \mathbf{D}^{\alpha+k-1} \circ \dots \circ \mathbf{D}^\alpha (\mathbf{L}_\alpha)^{-k/2}, \quad \text{alternatively} \quad (\mathbf{D}^\alpha)^k (\mathbf{L}_\alpha)^{-k/2}, \quad \alpha > -1.$$

If V is the operator defined by $Vf(y) = (2y)^{1/2}f(y^2)$ and $2\gamma = \eta + \frac{p}{2} - 1$, then $\|Vf\|_{L^p(y^\eta dy)} = 2^{\frac{1}{2} - \frac{1}{p}} \|f\|_{L^p(y^\gamma dy)}$ and the following proposition holds.

Proposition 6. *Let $1 < p < \infty$, $2\gamma = \eta + \frac{p}{2} - 1$ and T an operator defined over the set of finite linear combination of Laguerre functions $\{\mathfrak{L}_k^\alpha\}_k$. The operator T has a bounded extension from $L^p((0, \infty), y^\gamma dy)$ into $L^p((0, \infty), y^\gamma dy)$ if and only if the operator $\mathbf{T} = VT V^{-1}$ has a bounded extension from $L^p((0, \infty), y^\eta dy)$ into $L^p((0, \infty), y^\eta dy)$.*

An easy consequence of the above proposition, and Theorems 7 and 9 is the following.

Theorem 11. *Let $\alpha > -1$ and let f be a finite linear combination of Laguerre functions $\{\mathfrak{L}_k^\alpha\}_k$.*

- i) $e^{-tL_\alpha} f = V^{-1} e^{-t\mathbf{L}_\alpha} V f,$
- ii) $(L_\alpha)^{-s} f = V^{-1} (\mathbf{L}_\alpha)^{-s} V f,$ for all $s > 0,$
- iii) $\delta^\alpha f = V^{-1} \mathbf{D}^\alpha V f,$
- iv) $R_\alpha^k f = V^{-1} \mathbf{R}_\alpha^k V f.$

Proposition 7. *Let $\alpha > -1, 1 < p < \infty,$ and η be real number. Let S be any one of the operators $\mathbf{L}^{-s}, s > 0, \mathbf{R}_\alpha^k, (\mathbf{D}^\alpha)^k \mathbf{L}^{-k/2}, s > 0.$ Then the operator \mathbf{S} has a bounded extension from $L^p((0, \infty), y^\eta dy)$ into itself, for η satisfying*

$$-1 - \alpha p - \frac{p}{2} < \eta < \alpha p + \frac{3p}{2} - 1. \tag{16}$$

Now in the same way as in Section 3, we can define potential spaces and Sobolev spaces for the class of Laguerre functions $\{\varphi_k^\alpha\}_k, \alpha > -1.$ Thus, given $\alpha > -1, 1 < p < \infty, s > 0$ and η satisfying (16), we define

$$\mathfrak{U}_{\alpha,s}^p(y^\eta) = (\mathbf{L}_\alpha)^{-s/2} [L^p(\mathbb{R}^+, y^\eta dy)]$$

with the norm $\|f\|_{\mathfrak{U}_{\alpha,s}^p(y^\eta)} = \|g\|_{p,\eta},$ where $(\mathbf{L}_\alpha)^{-a/2} g = f.$

We shall denote by $U_\alpha^{k,p}(y^\gamma),$ the set of functions f in $L^p(\mathbb{R}^+, y^\eta dy)$ such that

$$\mathbf{D}^{\alpha+m} \circ \dots \circ \mathbf{D}^{\alpha+1} \circ \mathbf{D}^\alpha f \in L^p(\mathbb{R}^+, y^\eta dy), \quad 0 \leq m \leq k - 1,$$

with the norm

$$\|f\|_{U_{\alpha}^{k,p}(y^{\eta})} = \|f\|_{p,\eta} + \sum_{m=0}^{k-1} \|\mathbf{D}^{\alpha+m} \circ \dots \circ \mathbf{D}^{\alpha+1} \circ \mathbf{D}^{\alpha} f\|_{p,\eta}.$$

Finally, $\mathcal{U}_{\alpha}^{k,p}(y^{\eta})$, will denote the set of functions f in $L^p(\mathbb{R}^+, y^{\eta} dy)$ such that $(\mathbf{D}^{\alpha})^m f \in L^p(\mathbb{R}^+, y^{\eta} dy)$, $0 \leq m \leq k$, with the norm

$$\|f\|_{\mathcal{U}_{\alpha}^{k,p}(y^{\eta})} = \sum_{m=0}^k \|(\mathbf{D}^{\alpha})^m f\|_{p,\eta}.$$

The following theorems are direct consequences of Theorems 6, 10 and Propositions 7 and 11.

Theorem 12. *Let $\alpha > -1$, $1 < p < \infty$, $k \in \mathbb{N}$ and η satisfies (16).*

- i) $U_{\alpha,\eta}^{k,p} = \mathfrak{U}_{\alpha,\eta}^{k,p}$, and the norms are equivalent.
- ii) Let η satisfying $-\frac{\alpha}{2}p - 1 < \eta < p - 1 + \frac{\alpha}{2}p$. Then $\mathfrak{U}_{\alpha}^{k,p}(y^{\eta}) \subset \mathcal{U}_{\alpha}^{k,p}(y^{\eta})$.
- iii) Let $-1 < \alpha \leq 0$. Then $\mathfrak{U}_{\alpha}^{2,2} \neq \mathcal{U}_{\alpha}^{2,2}$.
- iv) If η satisfies

$$-1 - (\alpha - 1)p - \frac{p}{2} < \eta < (\alpha - 1)p + \frac{3p}{2} - 1, \tag{17}$$

then $\mathfrak{U}_{\alpha}^{2,p}(y^{\eta}) = \mathcal{U}_{\alpha}^{2,p}(y^{\eta})$.

Analogous results could be obtained for the systems of Laguerre functions $\ell_k^{\alpha}(y) = \mathfrak{L}_k^{\alpha}(y)y^{-\alpha/2}$ and $\psi_k^{\alpha}(y) = \sqrt{2}y^{-\alpha}\mathcal{L}_k^{\alpha}(y^2)$, $\alpha > -1$. These systems are eigenfunctions of the differential operators

$$\mathbb{L}_{\alpha} = -y \frac{d^2}{dy^2} - (\alpha + 1) \frac{d}{dy} + \frac{y}{4}.$$

and

$$\mathfrak{L}_{\alpha} = -\frac{1}{4} \left\{ \frac{d^2}{dy^2} + \left(\frac{2\alpha + 1}{y} \right) \frac{d}{dy} - y^2 \right\}.$$

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